

TOP QUARK  
FLAVOR PHYSICS ISSUES  
AT 8500'

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- $|V_{tb}|$
- Rare decays
- ~~CP~~

## TOP FLAVOR PHYSICS:

$|V_{tb}| = 0.9990$  to  $0.9993$  in the Standard Model.

This constraint is ENTIRELY from unitarity.

Our only "direct" measurement of  $|V_{tb}|$  is from

$$\text{BR}(t \rightarrow Wb) \approx 0.7 \text{ at CDF/D0.}$$

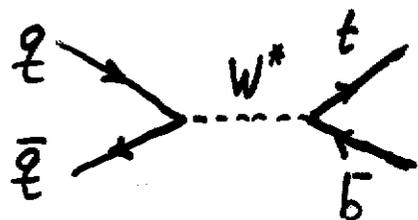
$\Rightarrow$  All this tells us is  $|V_{tb}|^2 \gg |V_{ts}|^2 + |V_{td}|^2$ .

WITHOUT UNITARITY:

$|V_{tb}|$  is essentially unconstrained.

Only direct probe of  $|V_{tb}|$  is  
 Single top production.

At hadron colliders, cleanest probe is "s-channel:"

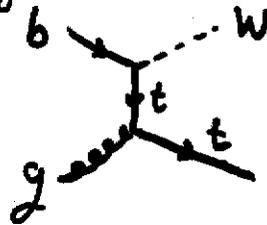
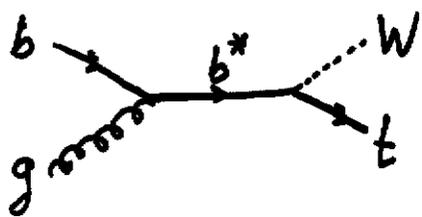


882 fb (TeV Run II)  
 10215 fb (LHC) ← higher backgr.

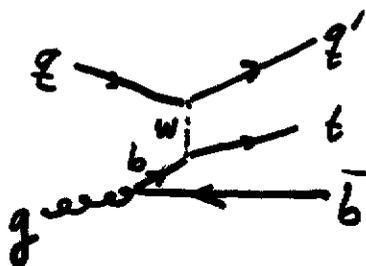
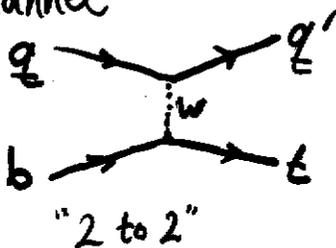
Given acceptances, backgrounds, we can expect a  
 few hundred events at Run 2a  $\Rightarrow$   $\sim 10\%$  cross-sec. measurement

Other diagrams contribute too — but final states are distinguishable:

"W-t associated production" of beauty sea — depends on  $b_g$  pdf's:



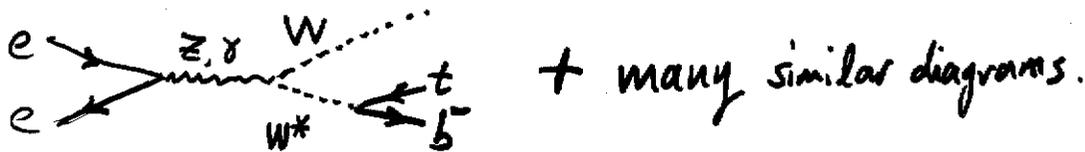
t-channel:



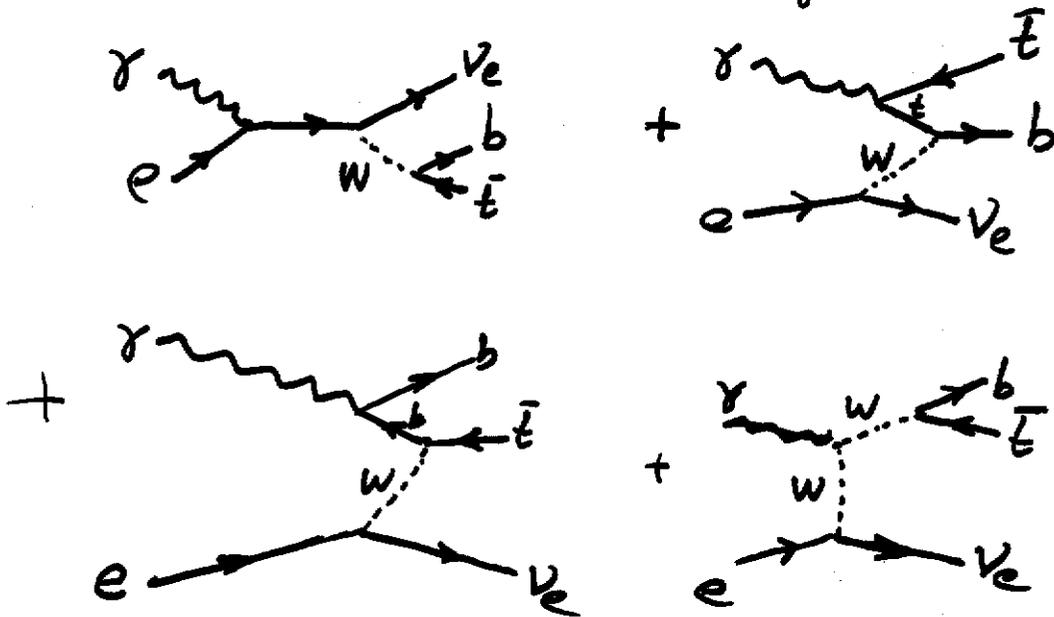
These processes  
 dominate the  
 s-channel at  
 LHC —  $|V_{tb}|$  for  
 these will be  
 QCD-limited.

ELECTRON COLLIDERS offer a theoretically cleaner probe of the W-t-b vertex:

$e^+e^- \rightarrow$  Cross-sections low ( $\ll 1$  fb)  
 same final state as  $e^+e^- \rightarrow t\bar{t}$



$e^-\gamma \rightarrow$  Much more promising!



$\sigma \approx 40$  fb ( $\sqrt{s} = 400$  GeV, unpolarized)  
 doubles for polarized beams.

$\approx 1\%$  measurement of  $|V_{tb}|$  possible!

# New physics observable in single top production:

DEFINE  $R_t \equiv \frac{\sigma(\text{single top})_{\text{observed}}}{\sigma(\text{single top})_{\text{SM}, V_{cb} \approx 1}}$

(Note that  $R_t$  could be different for different production channels.)

- 4<sup>th</sup> quark generation:

$$R_t < 1 \quad (|V_{cb}| < 1)$$

(Different 1-top channels give same  $R_t$ )

- Heavy gauge bosons ( $W'$ ) a la Technicolor:

$$R_t > 1 \quad (|V_{cb}| \approx 1)$$

(Different 1-top channels are at different  $Q^2 \Rightarrow$  different  $R_t$ )

Sensitivity at TeV to  $M_{W'} \sim 1-3$  TeV depending on mixing.

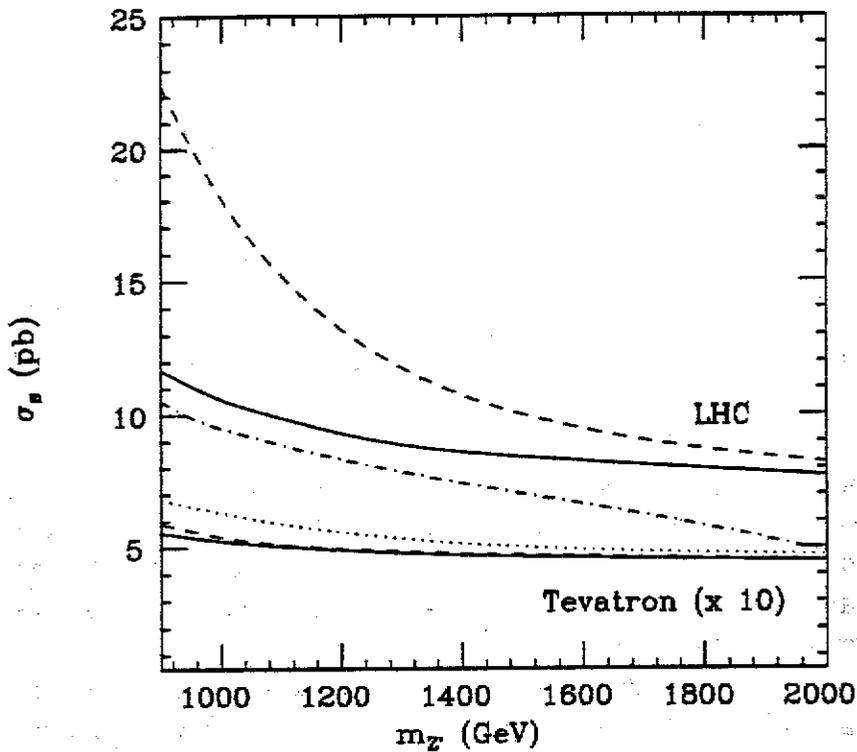


FIG. 7. The NLO rate of  $q \bar{q}' \rightarrow W, W' \rightarrow t \bar{b}$  ( $\sigma_S$ ) in pb at the Tevatron (lower curves) and CERN LHC (upper curves), for the top-quark-flavor model with  $\sin^2 \phi = 0.05$  (solid curves) and  $\sin^2 \phi = 0.25$  (dashed curves), as a function of  $M_{Z'} = M_{W'}$ . The Tevatron cross sections are multiplied by a factor of 10. At the Tevatron, the  $\bar{t}$  production rate is equal to the  $t$  rate. At the CERN LHC the  $\bar{t}$  rates are shown for  $\sin^2 \phi = 0.05$  (dotted curve) and  $\sin^2 \phi = 0.25$  (dot-dashed curve).

# RARE TOP DECAYS - $W+q$

In the SM, rare decays probably not observable:

$$\begin{array}{ll} t \rightarrow Ws & 1.6 \times 10^{-3} \\ t \rightarrow Wd & 10^{-4} \end{array}$$

Certainly produced copiously at TeV, LHC but background from mistagged  $b$  jets prohibitive.

- 4th Generation: Unitarity constraints on  $\Gamma(t \rightarrow Wb)$  disappear  $\Rightarrow$  only limits on  $\left( \frac{\Gamma(t \rightarrow Ws, Wd)}{\Gamma(t \rightarrow Wb)} \right)$  are from direct searches (CDF, D0).

$$BR(t \rightarrow Ws) < 0.25$$

$$BR(t \rightarrow Wd) < 0.01 \quad (\text{from loop-process limits on } |V_{td}|)$$

$b'$  may be light enough to show up in  $t$  decay:

$$BR(t \rightarrow W \underset{b' \text{ or } c}{b'}) \sim 10^{-3} \quad \text{for } m_{b'} \sim 100 \text{ GeV}/c^2$$

# RARE TOP DECAYS - FCNC

$$t \rightarrow c g$$

$$t \rightarrow c \gamma$$

$$t \rightarrow c Z$$

Only proceed at loop level in SM -  
 $BR \sim 10^{-13} - 10^{-12}$ .

THESE DECAYS ARE ENHANCED IN NEW PHYSICS MODELS:

- MSSM: All these decays factor of  $\sim 10^5$  above SM.  
 $\Rightarrow$  STILL not observable at TeV, LHC

- MSSM  $\mathcal{R}$ : Decays increase  $10^8 - 10^9$  above SM.  
 $\Rightarrow t \rightarrow c \gamma$  may be observable at LHC; others within 1 order of magnitude.

- Multiple Higgs Doublets: "Type III" models - Decays increase by factor of  $10^6 - 10^7$  above SM. Not observable soon.

# RARE TOP $\rightarrow$ HIGGS DECAYS:

## DECAYS TO NEUTRAL HIGGS:

SM: •  $BR(t \rightarrow Wbh^0) \lesssim 7 \times 10^{-8}$  for  $m_H \geq 100 \text{ GeV}/c^2$  IN SM  
→ DROPS RAPIDLY FOR HIGHER  $m_H$ ; W or  $h^0$  OFF-SHELL ALREADY.

•  $BR(t \rightarrow ch^0) = 10^{-12} \sim 10^{-14}$  for  $m_H = 100 \sim 160 \text{ GeV}/c^2$  in SM.

MSSM: •  $BR(t \rightarrow ch^0)$  UP TO  $10^{-4}$

2 HIGGS •  $BR(t \rightarrow ch^0)$  UP TO  $10^{-2}$  IN MODEL III  
(due to tree-level  $t\bar{t}H$  coupling  $\propto \sqrt{m_b m_c}$ )

## DECAYS TO CHARGED HIGGS:

MSSM: •  $BR(t \rightarrow h^+ b)$  potentially large (0.1) for large or small  $\tan\beta$ ;  
→ perhaps observable at TeV.

2 HIGGS: •  $BR(t \rightarrow h^+ b)$  potentially large

# RARE $t$ DECAY MATRIX

Decay	Predicted Rates								Sensitivity		
	SM	MSSM RPC	MSSM RPV	4th Gen	2 Higgs Doubl's	Ext. Dim.	Anom. Coupl.	Topcolor Asst. TC	TeV Run 2	LHC <sup>b</sup>	LC
$Wb$	$> 0.99$									$0.002^{p,C}$	
$W_s$	$1.6 \cdot 10^{-3} B$			$< 0.25^B$							
$Wd$	$1.0 \cdot 10^{-4} B$			$< 0.01^B$							
$h^+b$	0	$< 1^c$	$< 1^c$		$\leq BR(Wb)$				$11\%^a$	$0.03^{d,C}$	?
$cg$	$5 \cdot 10^{-11} B$	$10^{-6}$	$10^{-3} e$		$10^{-5} f,B$					$.0074^{g,C}$	
$c\gamma$	$5 \cdot 10^{-13} B$	$10^{-8}$	$10^{-4} e$		$10^{-7} f,B$				$\sim 0.1^b$	$.0001^{g,C}$	
$cZ$	$1.3 \cdot 10^{-13} B$	$10^{-8}$	$10^{-5} e$		$10^{-8} f,B$					$.0002^{g,C}$	
$ch^0$	$< 10^{-7} A$	$\sim 10^{-4} i$	$\sim 10^{-5} e,j$		$< 0.01$						
$ch^0$	$10^{-12} k$										
$ch^0$	$10^{-15} l$										
$WbZ$	$\sim 10^{-6} B$				$0.01^{f,m,B}$					$10^{-3} C$	
$Wbh^0$	$7 \cdot 10^{-8} k,B$										
$cWW$	$10^{-13} B$										
$Wb'$	0			$10^{-3} n,B$							

<sup>a</sup>Integrated luminosity  $2 \text{ fb}^{-1}$ .

<sup>b</sup>Integrated luminosity  $10 \text{ fb}^{-1}$ .

<sup>c</sup>Can be large for large or small values of  $\tan \beta$ .

<sup>d</sup>Limited by systematic uncertainty in  $\tau$  ID eff and fake  $t$ 's.

<sup>e</sup>In models with baryon number violating couplings.

<sup>f</sup>In Type III models (FCNC couplings allowed)

<sup>g</sup>Integrated luminosity  $100 \text{ fb}^{-1}$ .

<sup>i</sup>For  $CP$ -even higgs  $h^0$ .

<sup>j</sup>Using 2s bounds on relevant  $R$ -parity violating couplings. (??)

<sup>k</sup> $M_h = 100 \text{ GeV}/c^2$

<sup>l</sup> $M_h = 160 \text{ GeV}/c^2$

<sup>m</sup>Extension of 2HDM model to include Higgs triplets

<sup>n</sup> $M_{b'}$  =  $100 \text{ GeV}/c^2$

<sup>p</sup>Statistical sensitivity only; systematics may dominate.

<sup>A</sup>hep-ph/9802305.

<sup>B</sup>B.Mele, hep-ph/0003064.

<sup>C</sup>Atlas Detector and Physics Performance Technical Design Report (1999).

$\not{CP}$  in  $t \rightarrow Wb$  decay

Two-Higgs doublet models predict a CP-violating polarization of the W in  $t \rightarrow Wb$  decay.

$\Rightarrow$  Manifested experimentally as an asymmetry  $A^t \equiv \frac{N^{up} - N^{down}}{N^{up} + N^{down}}$  where  $N^{up}, N^{down}$  are the ~~the~~ number of W-decay leptons emitted, respectively, above and below the  $t \rightarrow Wb$  decay plane.

$A^t$  has no tree-level contribution in SM

$\Rightarrow$  vanishingly small.

In favored 2HDM regions with  $\tan\beta \ll 1$ ,

$$|A^t| = 10^{-5} \sim 10^{-3}$$

# CONCLUSIONS—

- RICH PROGRAMS OF FLAVOR PHYSICS WITH THE TOP QUARK WILL BE POSSIBLE AT NEW FACILITIES, INCLUDING A LINEAR COLLIDER (ESPECIALLY ONE WITH  $e^+e^-$  COLLISIONS).
- 4<sup>th</sup> GENERATION SEARCHES LOOK MOST PROMISING IN SINGLE TOP PRODUCTION AND  $W_s, W_d$  DECAYS
- SUSY, 2-HIGGS-DOUBLET PHYSICS COULD APPEAR IN  $t \rightarrow$  HIGGS DECAYS
- FCNC DECAYS WITHOUT HIGGS IN FINAL STATE LOOK VERY DIFFICULT EVEN WITH NEW PHYSICS.
- $CP$  IN 2HDMs INTERESTING WITH  $>10^6$   $t$  DECAYS.